

Rational solutions of integrable nonlinear wave models

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HISTORY 1

LINEAR differential equations with constant coefficients
do not have rational solutions

(1977–78) Adler, Airault, McKean, Moser, Ablowitz, Newell, Satsuma
Korteweg-deVries equation $u_t + u_{xxx} - 6uu_x = 0$

$$u_n(x, t) = -2\partial_x^2 \log(P_n(x, t)) , \quad n \geq 0$$

Adler-Moser polynomials : $P_0 = 1 , P_1 = x , P_2 = x^3 + 12t , \dots$

$$u_0 = 0 , \quad u_1 = \frac{2}{x^2} , \quad u_2 = 6x \frac{x^3 - 24t}{(x^3 + 12t)^2} , \dots$$

Boussinesq equation $u_{tt} \pm u_{xxxx} + (u^2)_{xx} = 0$
motion of poles as many-body system

HISTORY 2

connection to Painlevé II and IV : (1959–1965) Yablonskii–Vorob’ev polynomials, (1999) Noumi, Yamada (generalized Hermite polynomials and generalized Okamoto polynomials)

defocusing Nonlinear Schroedinger equation $iu_t + u_{xx} - 2|u|^2u = 0$
 (1985) Nakamura, Hirota, (1996) Hone, (2006) Clarkson

$$u_n = \frac{g_n}{f_n} , \quad n \geq 0$$

focusing Nonlinear Schroedinger equation $iu_t + u_{xx} + 2|u|^2u = 0$
 (1983) Peregrine, (2010) Clarkson, Matveev

$$u_n = \frac{G_n}{F_n} e^{2it}, \quad n \geq 0$$

$$G_0 = 1, F_0 = 1, G_1 = 4x^2 + 16t^2 - 4it - 3, F_1 = 4x^2 + 16t^2 + 1, \dots$$

PEREGRINE LUMP

rational soliton as ratio of polynomials of degree 2

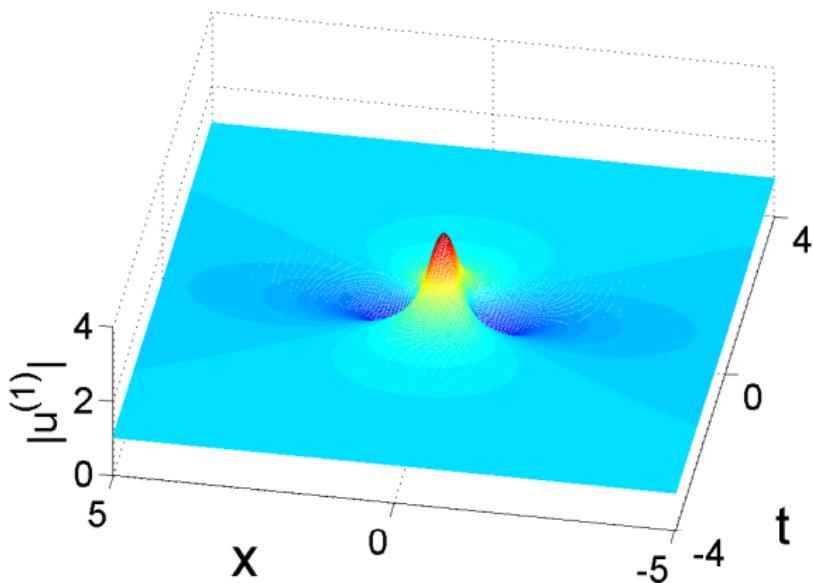


Figure: background amplitude=1 , peak amplitude = 3

HISTORY 3

" The finite density boundary conditions have meaningful applications only when $\chi > 0$, hence we shall confine ourselves to this case. "

L. Faddeev and L. Takhtajan *Hamiltonians Methods in the Theory of Solitons*, Springer (1986)

recent extensions to other integrable models such as:

- vector nonlinear Schroedinger equations
- Hirota equation and coupled Hirota equations
- three wave resonant interaction model
- Massive Thirring Model
- discrete NLS equation
- several others

GENERAL OBSERVATIONS

- making a limit :

$$M(z) = \sum_{j=1}^{N+1} \gamma_j e^{k_j z} \rightarrow e^{k_c z} P_{(N)}(z) = e^{k_c z} \sum_{j=0}^N c_j z^j, \quad k_j \rightarrow k_c$$

- computing the critical value k_c

Example : KdV for Adler-Moser polynomials , $k_c = 0$

Example : NLS for Peregrine and higher order , $k_c = \pm i$

NLS equation

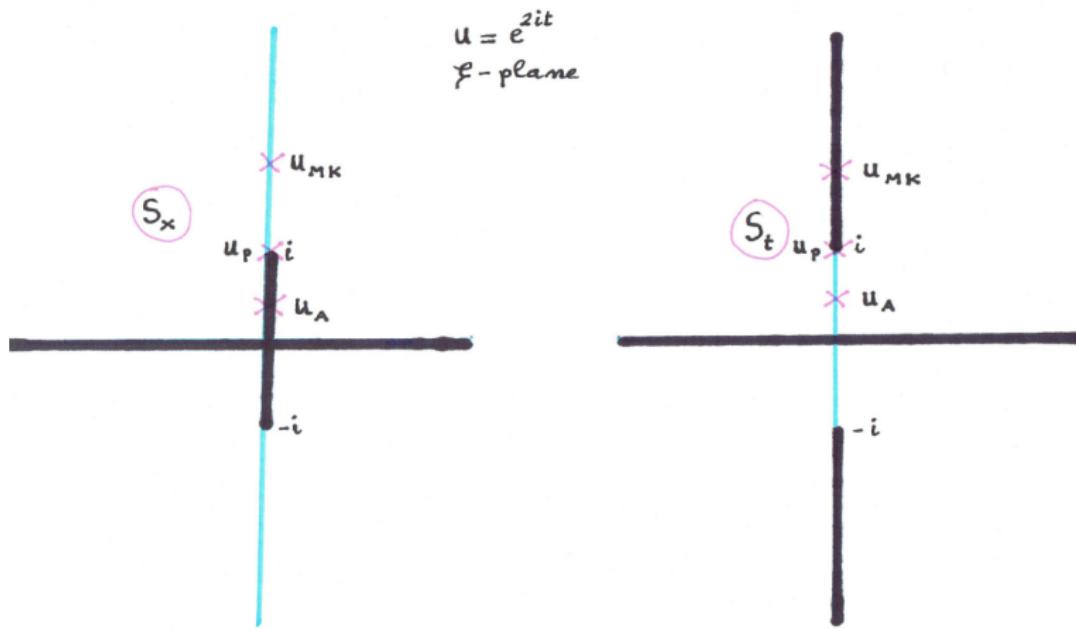


Figure: $S_x = x$ -part continuum spectrum / $S_t = t$ -part continuum spectrum

computing k_c 1

preliminary note on Jordan forms : $M = TM^{(J)}T^{-1}$

$$M^{(J)} = \{n_j \times n_j \text{ blocks}\} = \{m_j \mathfrak{J}_{n_j \times n_j} + \mu_j \mathfrak{J}_{n_j \times n_j}\}$$

$\mathfrak{J}_{n_j \times n_j}$ is the $n_j \times n_j$ unit matrix and $\mathfrak{J}_{n_j \times n_j} = \begin{pmatrix} 0 & 1 & \cdots & \cdots & 0 \\ 0 & 0 & 1 & \cdots & 0 \\ \vdots & \vdots & \vdots & \vdots & 1 \\ 0 & \cdots & \cdots & \cdots & 0 \end{pmatrix}$

n_j is the algebraic multiplicity of the eigenvalue m_j and $\mathfrak{J}_{n_j \times n_j}^{n_j} = 0$
if $N^n \neq 0$ and $N^{n+1} = 0$ then $e^{zN} = P_n(z)$

$$e^{zM} = T\{e^{zm_j} P_{n_j-1}(z)\} T^{-1}$$

necessary condition for $\mu_j \neq 0$ is $n_j > 1$

example : NLS equation

$$u_t = i[u_{xx} - 2s|u|^2] , \quad \Psi_x = X\Psi , \quad \Psi_t = T\Psi , \quad s = \pm 1$$

$$u_0(x, t) = ae^{-isa^2t} , \quad \Psi_0(x, t, k) = G(x, t)e^{i(\Lambda(k)x - \Omega(k)t)}$$

DEFINITION : k_c is a *critical value* of k if $\Lambda(k_c)$ is similar to a Jordan form Λ_J :

$$\Lambda(k_c) = T \Lambda_J T^{-1}$$

$$\Lambda(k) = \begin{pmatrix} k & -isa \\ -ia & -k \end{pmatrix} , \quad \lambda_1 = \sqrt{k^2 - sa^2} , \quad \lambda_2 = -\sqrt{k^2 - sa^2}$$

for $s = 1$, $k_c = \pm a$, for $s = -1$, $k_c = \pm ia$, $\Lambda^2(k_c) = 0$

$$e^{i\Lambda(k_c)x} = 1 + i\Lambda(k_c)x$$

study case : vector NLS equation 1)

$$\begin{cases} u_t^{(1)} = i[u_{xx}^{(1)} - 2(s_1 |u^{(1)}|^2 + s_2 |u^{(2)}|^2)u^{(1)}] \\ u_t^{(2)} = i[u_{xx}^{(2)} - 2(s_1 |u^{(1)}|^2 + s_2 |u^{(2)}|^2)u^{(2)}] \end{cases}$$

$$\Psi_x = X\Psi \quad , \quad \Psi_t = T\Psi$$

$$X(x, t, k) = ik\sigma + Q(x, t) \quad , \quad T = 2ik^2\sigma + 2kQ + i\sigma(Q^2 - Q_x)$$

$$\sigma = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix}, \quad Q = \begin{pmatrix} 0 & s_1 u^{(1)*} & s_2 u^{(2)*} \\ u^{(1)} & 0 & 0 \\ u^{(2)} & 0 & 0 \end{pmatrix}$$

study case : vector NLS equation 2)

$$\Psi(x, t, k) = \left[\mathbf{1} + \left(\frac{\chi - \chi^*}{k - \chi} \right) P(x, t) \right] \Psi_0(x, t, k)$$

$$\begin{pmatrix} u^{(1)}(x, t) \\ u^{(2)}(x, t) \end{pmatrix} = \begin{pmatrix} u_0^{(1)}(x, t) \\ u_0^{(2)}(x, t) \end{pmatrix} + \frac{2i(\chi - \chi^*)\zeta^*}{|\zeta|^2 - s_1|z_1|^2 - s_2|z_2|^2} \begin{pmatrix} z_1 \\ z_2 \end{pmatrix}$$

$$P(x, t) = \frac{ZZ^\dagger}{|\zeta|^2 - s_1|z_1|^2 - s_2|z_2|^2} \begin{pmatrix} 1 & 0 & 0 \\ 0 & -s_1 & 0 \\ 0 & 0 & -s_2 \end{pmatrix}$$

$$Z(x, t) = \begin{pmatrix} \zeta(x, t) \\ z_1(x, t) \\ z_2(x, t) \end{pmatrix} = \Psi_0(x, t, \chi^*) Z_0$$

study case : vector NLS equation 3)

$$\begin{pmatrix} u_0^{(1)}(x, t) \\ u_0^{(2)}(x, t) \end{pmatrix} = \begin{pmatrix} a_1 e^{i(qx - \nu t)} \\ a_2 e^{-i(qx + \nu t)} \end{pmatrix}, \quad \nu = q^2 + 2(s_1 a_1^2 + s_2 a_2^2), \quad a_j > 0$$

$$\Psi_0(x, t, k) = G(x, t) e^{i(\Lambda(k)x - \Omega(k)t)}, \quad [\Lambda(k), \Omega(k)] = 0$$

$$Z(x, t) = G(x, t) e^{i(\Lambda(\chi^*)x - \Omega(\chi^*)t)} Z_0$$

$$\Lambda(k) = \begin{pmatrix} k & -is_1 a_1 & -is_2 a_2 \\ -ia_1 & -k - q & 0 \\ -ia_2 & 0 & -k + q \end{pmatrix}$$

$$P_\Lambda(\lambda) = \det[\lambda - \Lambda(k)] = \lambda^3 + A_2(k)\lambda^2 + A_1(k)\lambda + A_0(k)$$

$$\Delta(k) = \text{discriminant of } P_\Lambda(\lambda) = k^4 + D_3 k^3 + D_2 k^2 + D_1 k + D_0$$

study case : vector NLS equation 4)

classification of rational solutions by computing :

- ① the critical value k_c

$$\Delta(k_c) = 0 , \quad k_c \neq k_c^*$$

- ② the similarity matrix T , the Jordan form Λ_J and the matrix $\widehat{\Omega}$

$$\Lambda(k_c) = T \Lambda_J T^{-1} , \quad \Omega(k_c) = T \widehat{\Omega} T^{-1} , \quad [\Lambda_J , \widehat{\Omega}] = 0$$

- ③ the vector

$$Z(x, t) = G(x, t) T e^{i(\Lambda_J x - \widehat{\Omega} t)} \begin{pmatrix} \gamma_1 \\ \gamma_2 \\ \gamma_3 \end{pmatrix}$$

CLASSIFICATION - 1

Case $[\lambda_1 = \lambda_2 = \lambda_3]$

$$q \neq 0, \quad a_1 = a_2 = 2q, \quad s_1 = s_2 = -1, \quad k_c = \pm i \frac{\sqrt{27}}{2} q$$

$$\Lambda_J = \begin{pmatrix} \lambda_1 & \mu_1 & 0 \\ 0 & \lambda_1 & \mu_1 \\ 0 & 0 & \lambda_1 \end{pmatrix}, \quad \widehat{\Omega} = \begin{pmatrix} \omega_1 & \rho_1 & \rho_2 \\ 0 & \omega_1 & \rho_1 \\ 0 & 0 & \omega_1 \end{pmatrix}$$

$$\lambda_1 = -\frac{k_c}{3}, \quad \mu_1 = 2iq, \quad \omega_1 = \frac{11}{2}q^2, \quad \rho_1 = 4\sqrt{3}q^2, \quad \rho_2 = 4q^2$$

$$T = \begin{pmatrix} \theta & 0 & -i \\ 1 & \theta^* & i\sqrt{3} \\ i\theta^* & i & 0 \end{pmatrix}, \quad \theta = \frac{1}{2}(-\sqrt{3} + i)$$

CLASSIFICATION - 2

1 $\gamma_3 = 0$

$$\begin{pmatrix} u^{(1)}(x, t) \\ u^{(2)}(x, t) \end{pmatrix} = \begin{pmatrix} e^{i(qx - \nu t)} & 0 \\ 0 & e^{-i(qx + \nu t)} \end{pmatrix} \frac{1}{P_2} \begin{pmatrix} P_2^{(1)} \\ P_2^{(2)} \end{pmatrix}$$

2 $\gamma_2 = 0$

$$\begin{pmatrix} u^{(1)}(x, t) \\ u^{(2)}(x, t) \end{pmatrix} = \begin{pmatrix} e^{i(qx - \nu t)} & 0 \\ 0 & e^{-i(qx + \nu t)} \end{pmatrix} \frac{1}{P_4} \begin{pmatrix} P_4^{(1)} \\ P_4^{(2)} \end{pmatrix}$$

VNLS rational solutions 1 ($\lambda_1 = \lambda_2 = \lambda_3$)

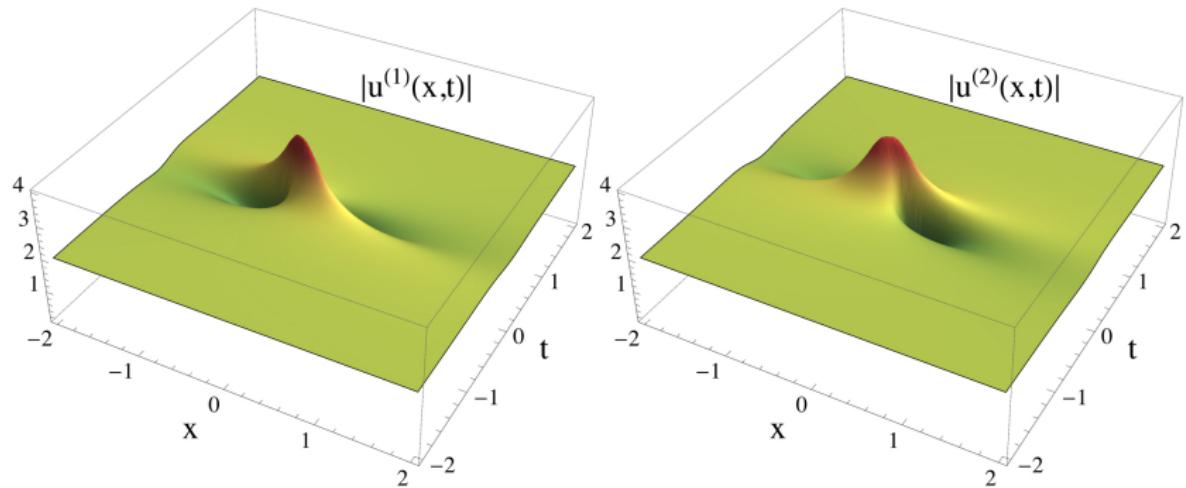


Figure: $k_c = i\frac{\sqrt{27}}{2}$, $s_1 = s_2 = -1$, $q = 1$, $a_1 = a_2 = 2$; $\gamma_2 = 1$, $\gamma_1 = \gamma_3 = 0$.

VNLS rational solutions 2 ($\lambda_1 = \lambda_2 = \lambda_3$)

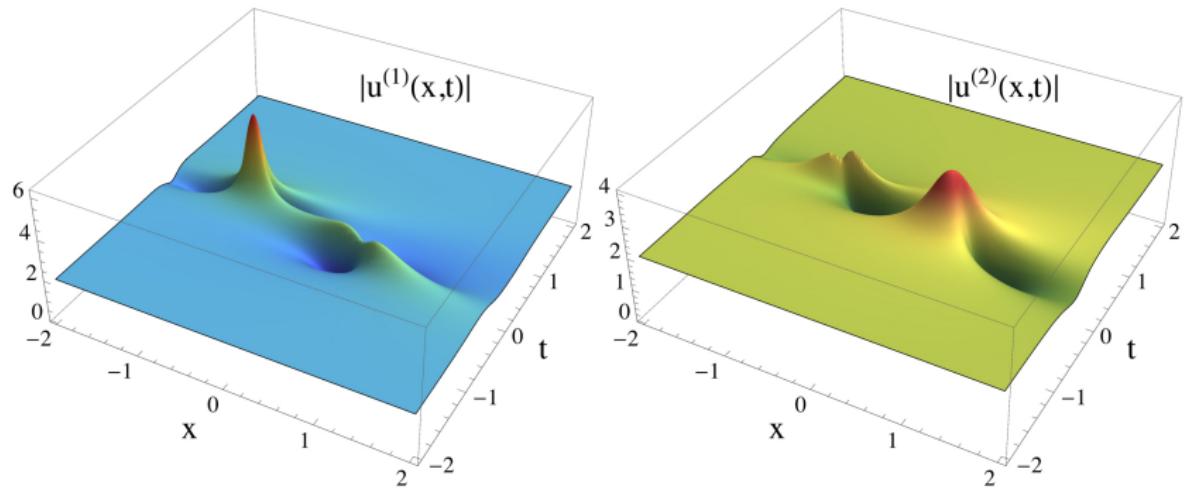


Figure: $k_c = i\frac{\sqrt{27}}{2}$, $s_1 = s_2 = -1$, $q = 1$, $a_1 = a_2 = 2$, $\gamma_1 = i$, $\gamma_2 = 0$, $\gamma_3 = 1$.

CLASSIFICATION - 3

Case $[\lambda_1 = \lambda_2 \neq \lambda_3]$

$$\Lambda_J = \begin{pmatrix} \lambda_1 & \mu & 0 \\ 0 & \lambda_1 & 0 \\ 0 & 0 & \lambda_3 \end{pmatrix}, \quad \hat{\Omega} = \begin{pmatrix} \omega_1 & \rho & 0 \\ 0 & \omega_1 & 0 \\ 0 & 0 & \omega_3 \end{pmatrix}$$

- ① $q = 0, s_1 = s_2 = -1$ explicit analytical
- ② $q \neq 0, s_1 = s_2, a_1 = a_2$ explicit analytical
- ③ $q \neq 0, a_1 \neq a_2$ numerical

VNLS rational solutions 3 ($\lambda_1 = \lambda_2 \neq \lambda_3$)

$q = 0$, $s_1 = s_2 = -1$ vector Peregrine solution

$$\begin{pmatrix} u^{(1)}(x, t) \\ u^{(2)}(x, t) \end{pmatrix} = e^{2i\omega t} \left[\frac{L}{B} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} + \frac{M}{B} \begin{pmatrix} a_2 \\ -a_1 \end{pmatrix} \right]$$

$$L = P_2 + |f|^2 e^{2px} , M = 4fe^{px+i\omega t}P_1 , B = \hat{P}_2 + |f|^2 e^{2px}$$

$$k_c = \pm ip , p = \sqrt{a_1^2 + a_2^2} , \omega = a_1^2 + a_2^2$$

$$\lambda_1 = \lambda_2 = 0 , \lambda_3 = -ip , \mu = -ip , \omega_1 = \omega_2 = p^2 , \omega_3 = 0 , \rho = -2p^2$$

$$T = \begin{pmatrix} -p & p & 0 \\ a_1 & 0 & a_2 \\ a_2 & 0 & -a_1 \end{pmatrix}$$

VNLS rational solutions 4 ($\lambda_1 = \lambda_2 \neq \lambda_3$)

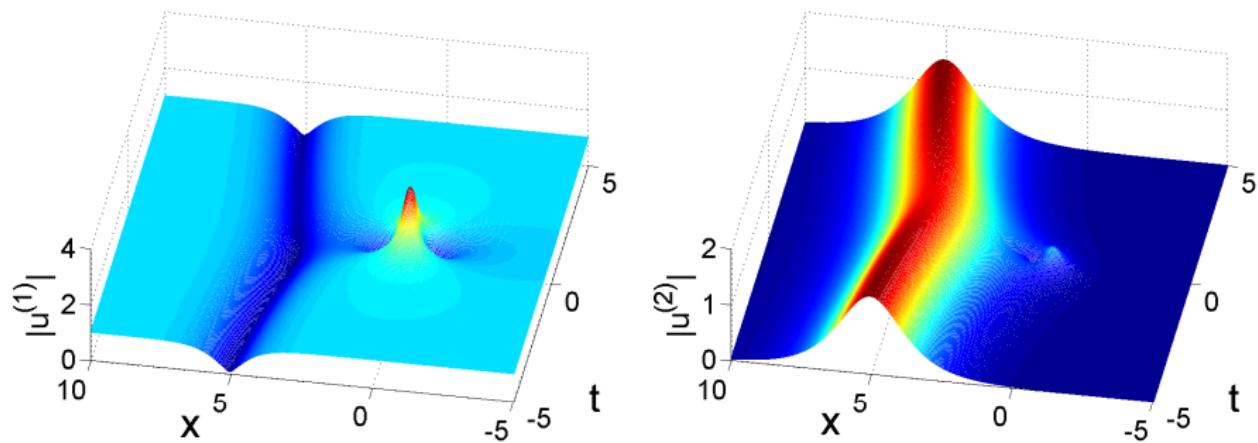


Figure: $k_c = i, q = 0, a_1 = 1, a_2 = 0, s_1 = s_2 = -1, f = 0.1,$

VNLS rational solutions 5 ($\lambda_1 = \lambda_2 \neq \lambda_3$)

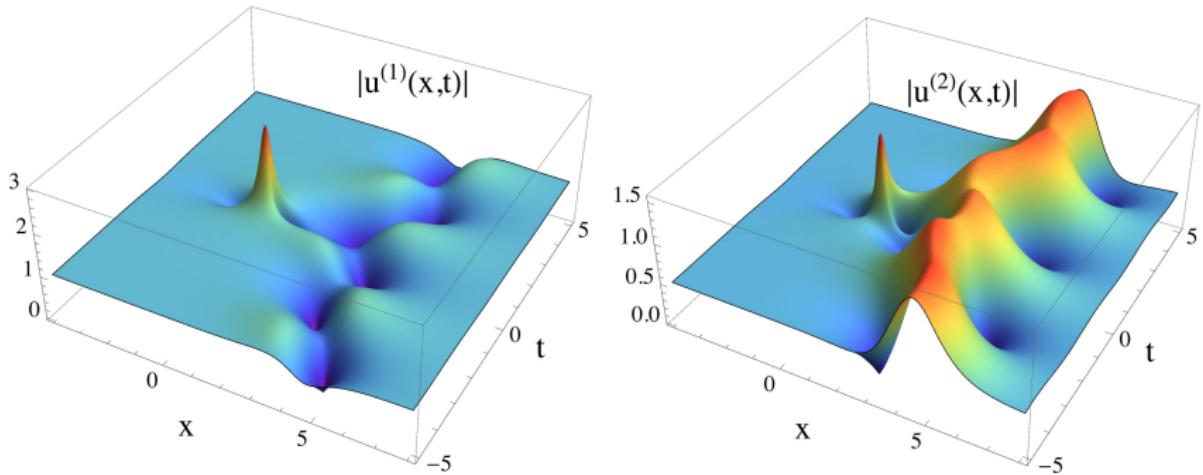


Figure: $k_c = i\frac{\sqrt{5}}{2}, q = 0, a_1 = 1, a_2 = 0.5, s_1 = s_2 = -1, f = 0.1i$

VNLS rational solutions 6 ($\lambda_1 = \lambda_2 \neq \lambda_3$)

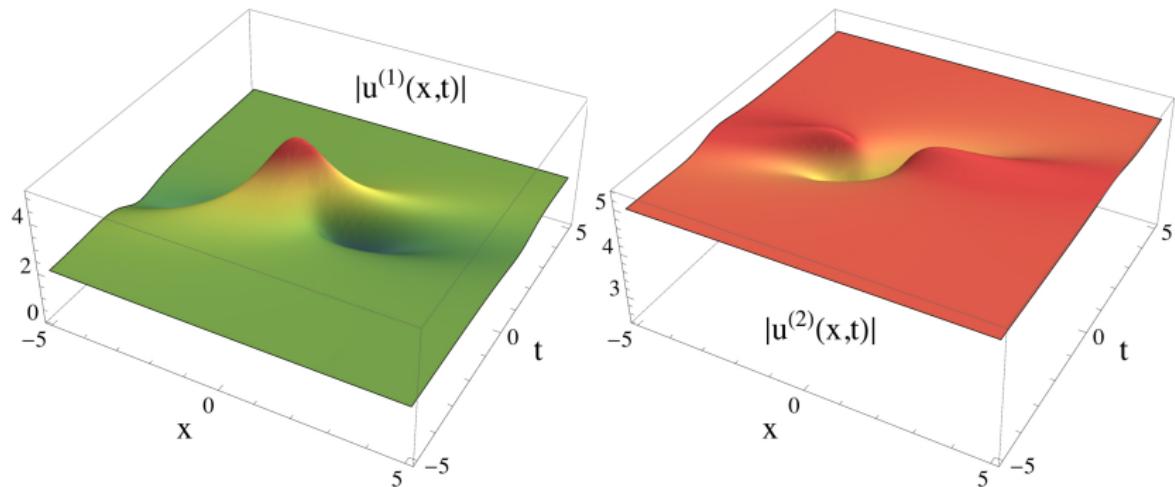


Figure:

$k_c = 4.876 + 5.343i$, $q = 1$, $a_1 = 2$, $a_2 = 5$, $s_1 = s_2 = -1$, $\gamma_2 = 1$, $\gamma_1 = \gamma_3 = 0$

VNLS rational solutions 7 ($\lambda_1 = \lambda_2 \neq \lambda_3$)

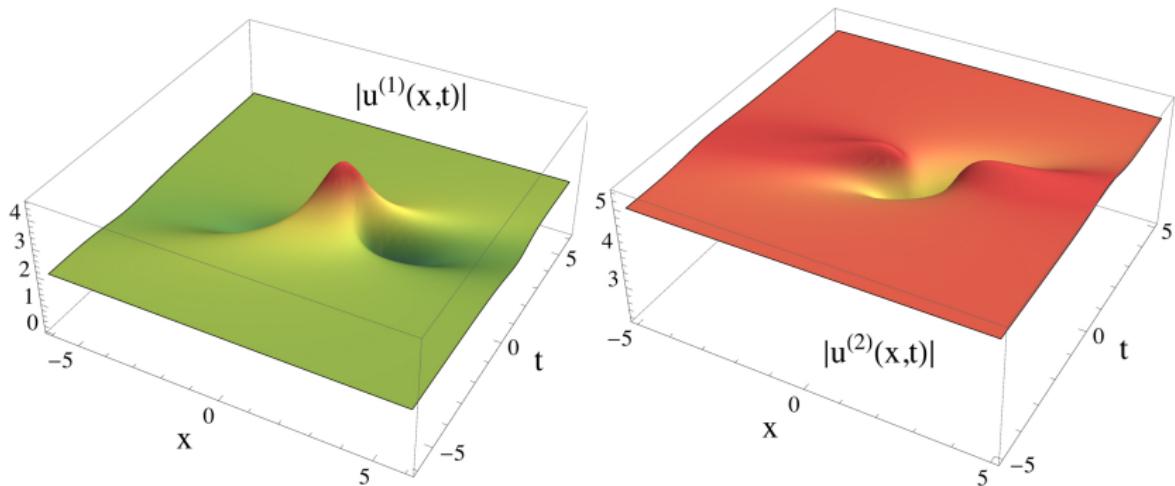


Figure:

$k_c = -5.600 + 4.655i$, $q = 1$, $a_1 = 2$, $a_2 = 5$, $s_1 = s_2 = 1$, $\gamma_2 = 1$, $\gamma_1 = \gamma_3 = 0$

VNLS rational solutions 8 ($\lambda_1 = \lambda_2 \neq \lambda_3$)

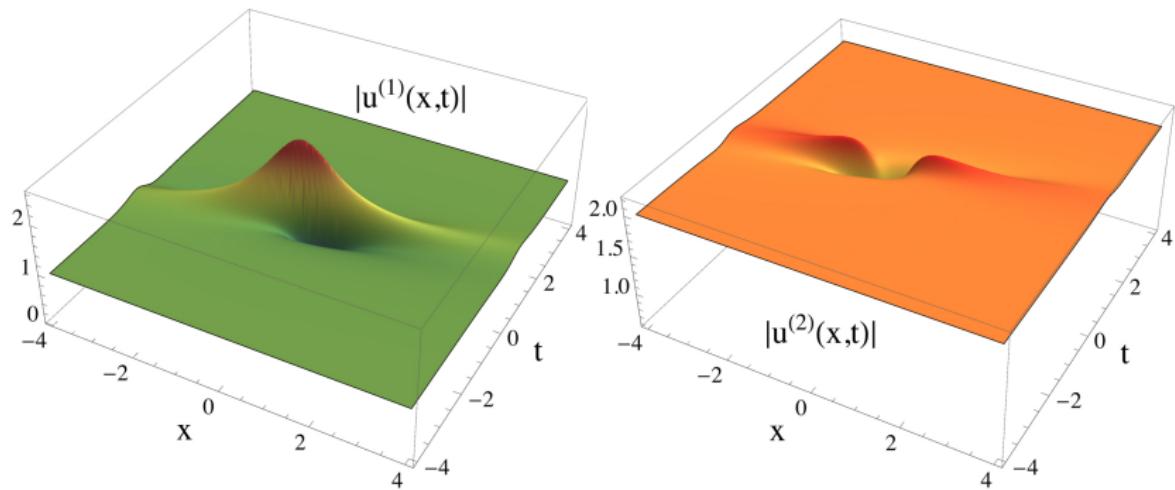


Figure: $k_c = -1.242 + 0.636i$, $q = 1$, $a_1 = 2$, $a_2 = 2$, $s_1 = -1$, $s_2 = 1$, $\gamma_2 = 1$, $\gamma_1 = \gamma_3 = 0$

other integrable equations

3 wave resonant interaction equations :

$$\begin{cases} E_{1t} + V_1 E_{1x} = E_2^* E_3^* \\ E_{2t} + V_2 E_{2x} = -E_1^* E_3^* \\ E_{3t} + V_3 E_{3x} = E_1^* E_2^* \end{cases}$$

Massive Thirring Model equations :

$$\begin{cases} iU_\xi - \nu V = \frac{1}{\nu} |V|^2 U \\ iV_\eta - \nu U = \frac{1}{\nu} |U|^2 V \end{cases}$$

$$\partial_\xi = \partial_t + c \partial_x , \quad \partial_\eta = \partial_t - c \partial_x$$

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